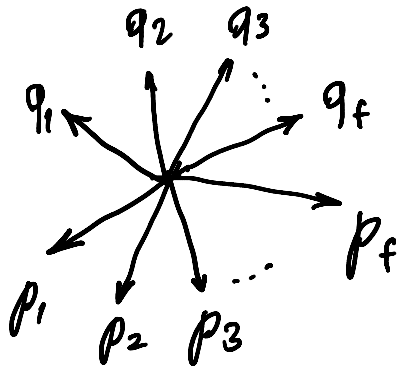
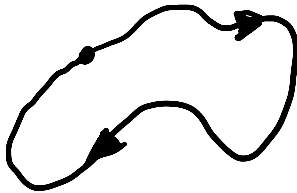


Phase space. Distribution function.

Phase space - space consisting of coordinates and momenta



Because Hamilton's equations are 1st-order differential equations, a point in phase space specifies uniquely all the subsequent evolution



Phase trajectory

As an example, consider a harmonic oscillator

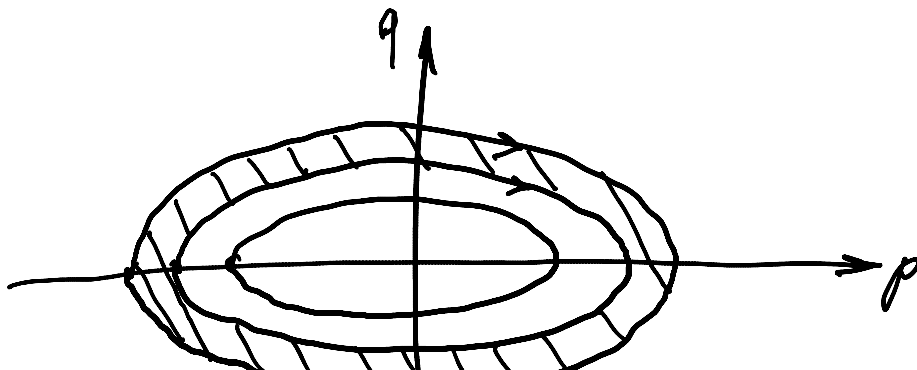
$$H = \frac{p^2}{2m} + \frac{m\omega^2 q^2}{2}$$

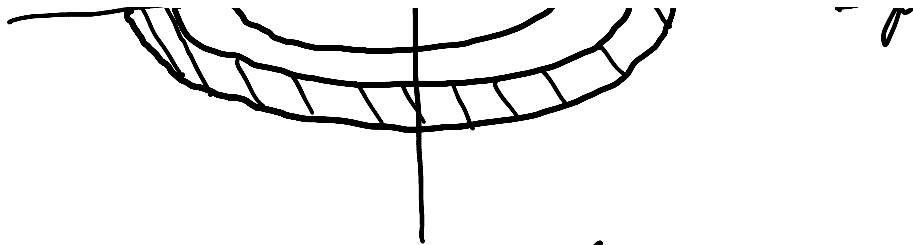
$$\begin{cases} \dot{p} = -m\omega^2 q \\ \dot{q} = \frac{p}{m} \end{cases} \rightarrow \ddot{q} + \omega^2 q = 0$$

$$q = A \sin(\omega t + \varphi)$$

$$p = Am\omega \cos(\omega t + \varphi)$$

$$\left(\frac{q}{A}\right)^2 + \left(\frac{p}{m\omega A}\right)^2 = 1 \quad \text{- equation of motion in phase space}$$





The area of the ellipse  $S = \pi m \omega A^2$

$$\text{Energy } \varepsilon = \frac{p^2}{2m} + \frac{m \omega^2 q^2}{2} = \frac{1}{2} A^2 m \omega^2 \cos^2 + \frac{1}{2} A^2 m \omega^2 \sin^2 = \frac{1}{2} A^2 m \omega^2$$

$$S = \oint p dq = 2\pi \frac{\varepsilon}{\omega} = 2\pi \hbar N$$

Transition to quantum case

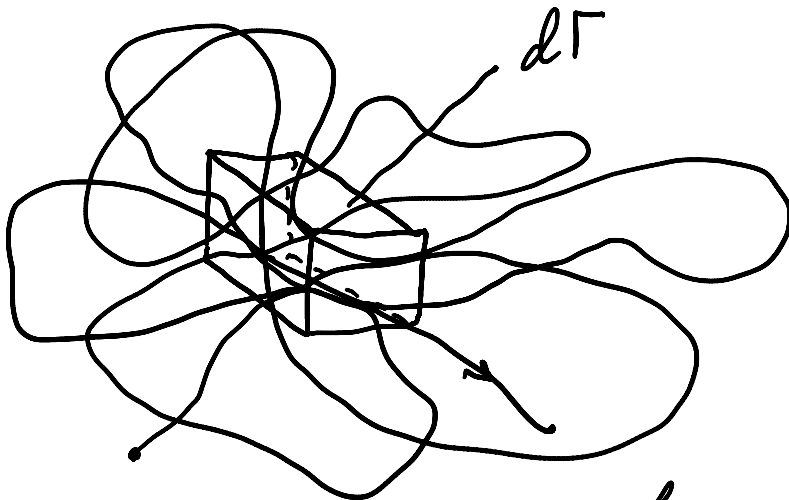
The concept of phase space is still very useful for a system of many particles

For  $N$  particles in 3D

$$d\Gamma = \underbrace{dq_1 dq_2 \dots dq_{3N}}_{\text{in the configuration space}} \underbrace{dp_1 dp_2 \dots dp_{3N}}_{\text{phase space } dp}$$

$dq$  Infinitesimal el-nt  
in the configuration space

phase space  $dp$



The statistical approach is based on the assumption that the system will go through a  $\Omega$  +  $\Gamma$  phase space many times

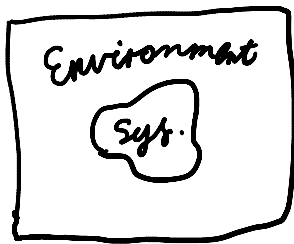
assumption that the system will visit  
 given element of phase space many times

$$dW = \lim_{T \rightarrow \infty} \frac{dt}{T} \quad \text{- prob-ty to find the system there}$$

$$dW = \underbrace{\rho(p_1, p_2, \dots, p_f; q_1, q_2, \dots, q_f)}_{\text{Statistical distribution function}} dp dq$$

$$\int \rho dp dq = 1$$

In an equilibrium system the distribution function is independent of the initial conditions. That applies to a system in an external environment as well.



That problem is way easier than solving the microscopic equations of motion

The probability of quantity  $f$ :

$$\bar{f} = \int p(p, q) \rho(p, q) dp dq$$

If there are 2 independent systems

$$\rho_{12} dp^{(12)} dq^{(12)} = \rho_1 dp^{(1)} dq^{(1)} \cdot \rho_2 dp^{(2)} dq^{(2)}$$

$$\rho_{12} = \rho_1 \rho_2$$

$$\langle f_1, f_2 \rangle = \langle f_1 \rangle \langle f_2 \rangle$$